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#### **Abstract**

The Boltzmann–Gibbs (BG) entropy and its associated statistical mechanics were generalized, three decades ago, on the basis of the nonadditive entropy ( ), which recovers the BG entropy in the limit.

The optimization of under appropriate simple constraints straightforwardly yields the so-called q-(https://www.cookiebot.com/en/what-exponential and q-Gaussian distributions, respectively generalizing the exponential and Gaussian ones, recovered for . These generalized functions ubiquitously emerge in complex systems, especially as economic and financial stylized features. These include price returns and volumes distributions, inter-occurrence times, characterization of wealth distributions and associated inequalities, among others. Here, we briefly review the basic concepts of this q-statistical

Keywords: economics and finance (/search?q=economics+and+finance); nonadditiwe entropies (/search?q=nonadditive+entropies); nonextensive statistical mechanics (/search?

generalization and focus on its rapidly growing applications in economics and finance.

q=nonextensive+statistical+mechanics)

### 1. Introduction

Nexponential and Gaussian functions ubiquitously emerge within linear theories in mathen...dics, physics, economics and elsewhere. To illustrate in what sense they are linear, let us focus on three typinal methods tituations, namely an ordinary differential equation, a partial derivative puation and an entropic optimization.

Consider the following ordinary differential equation: **Statistics** 

The solution is the well-known exponential function:

Consider now the following partial derivative equation:

where is the Dirac delta function. The solution is the well-known Gaussian distribution:

Let us finally consider the following entropic functional:

with the constraint:

where BG stands for Boltzmann–Gibbs; k is a conventional positive constant (usually in physics, and

elsewhere). If we optimize the functional (5) with the constraint (6) and:

being bounded below, we obtain:

(8)

(7)

(1)

(2)

(3)

(4)

(5)

(6)

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where is the Lagrange parameter associated with Constraint (7); T is the absolute temperature within BG statistical mechanics (necessarily if is unbounded from above; but both and possibilities exist if is bounded also from above). The probability distribution corresponds to the celebrated BG weight, where is usually referred to as the partition function. Two particular cases emerge frequently. The first

of them is with , hence , thus recovering solution (2). The second one is with , hence , thus recovering solution (4). Therefore, basic cases connect with the solu**(fictips:://t/wwwineaolFiebattocrsn(//e/a/wch/at)**. In addition to that, let us make explicit in what sense is itself**ishbahind-powertedaby/seekichet**ituted by two probabilistically independent subsystems A and B. In other words, we consider the case where

the joint probability of factorizes, i.e., . We straightforwardly verify that the functional is additive in the sense of Penrose [1], namely that: In the present brief review, we shall address a special class of nonlinearities, namely those

emerging within nonextensive statistical mechanics, q-statistics for short [2,3,4,5,6]. Equation (1) is now generalized into the following nonlinear one: (10)

Its solution is:

**Necessary** 

where the q-exponential function is defined as , with if and zero otherwise. Its inverse function is the q-logarithm, defined as . To avoid any confusion, let us mention that many other q-deformations of the

exponential and logarithmic functions have been introduced in the literature for a variety of purposes; among them, we have for instance Ramanujan's *q*-exponential function, unrelated to the present one.

Statistics (3) is now generalized into the following nonlinear one (referred to in the literature as

the porous medium equation [7,8,9]): Marketing

Its solution generalizes Equation (4) and is given by:

economic- and financial-specific features presented later on.

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validated (within a experimental error) in a granular medium [10]. The important connection between

the power-law nonlinear diffusion (12) and the entropy described here below was first established by Plastino and Plastino in [11], where they considered a more general evolution equation that reduces to (12) in the particular case of vanishing drift (i.e., ). The Plastino–Plastino Equation [11] with generalizes the porous medium equation in the same sense that the linear Fokker–Planck equation generalizes the classical heat equation. The above nonlinear Equations (10) and (12) have been

addressed here in order to provide some basic mathematical structure to approaches of various

Let us now focus on the entropic functional upon which nonextensive statistical mechanics is

Show details >

(12)

(11)

(13)

(14)

Before going on, let us mention that solution (13) implies that scales like, hence normal diffusion

for,anomalous sub-diffusion for and super-diffusion for,which has recently been impressively

with:

(16)

(18)

systems

with . If we optimize this functional with the constraints (6) and:

we obtain [**4**]:

hence

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As before, two particular cases emerge frequently. Tis-helingwar with y reactive to the form of (11). The second one is with ; hence, recovers the form of solution (13). Finally, if we consider itself for two independent subsystems A and B, we straightforwardly verify the following nonlinear composition law:

We then say that is nonadditive for . Entropic additivity is recovered if , which can occur in two different circumstances: for fixed k or for fixed q. Since k always appears in physics in the form, the limit is equivalent to . This is, by the way, the basic reason for which, in the limit of high temperatures or Newesenergies, Maxwell–Boltzmann statistics, Fermi–Dirac, Bose–Einstein and *q-*st*ɛ′* `tics

asymptotically coincide. The above q-generalized thermostatistical theory has been useful in the study of a considerable

systems

number of natural, artificial and social systems (see [**12**]). Theoretical and experimental illustracions in

include long-range-interacting many-body classical Hamiltonian

[**135t4,t15t16**s17,18,19,20] (see also [21,22]; the study of the long-range version of [23] wor be interesting), dissipative many-body systems [24], low-dimensional dissipative and conservative

nonlinear dynamical systems [25,26,27,28,29,30,31], cold atoms [32,33,34], plasmas [35,36] trapped atoms [37], spin-glasses [38], power-law anomalous diffusion [39,40], granular matter [1u], high-

earthquakes [50], biology [51,52], solar wind [53,54], anomalous diffusion in relation to central limit theorems and overdamped systems [55,56,57,58,59,60,61,62,63,64], quantum entangled systems [65,66], quantum chaos [67], astronomical systems [68,69], thermal conductance [70], mathematical

energy particle collisions [41,42,43,44,45,46], black holes and cosmology [47,48], chemistry [49],

structures [71,72,73,74,75,76] nonlinear quantum mechanics and

[77,78,79,80,81,82,83,84,85,86,87,88,89,90,91,92,93,94,95,96], others. Illustrations among in artificial systems include signal and image processing [97,98] and (asymptotically) scale-free

networks [99,100,101]. In the realm of social systems, from now on, we focus on economics and

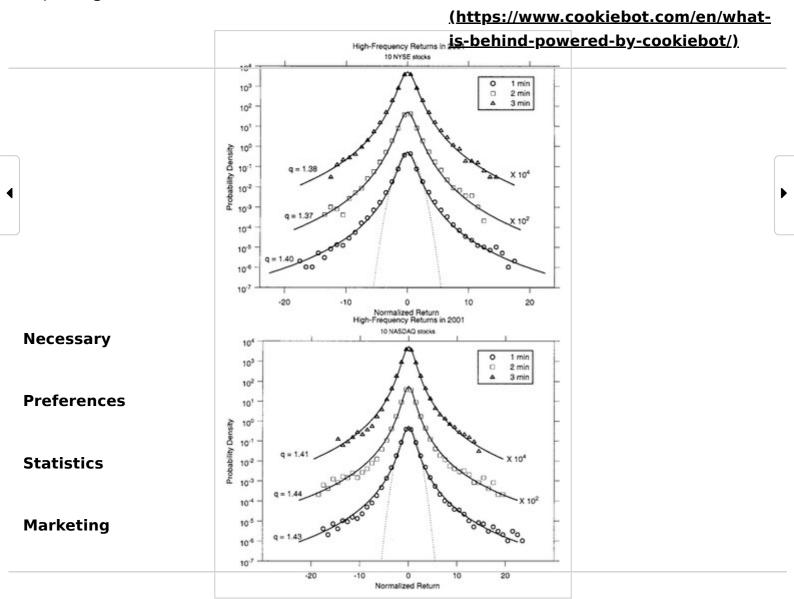
financial theory [102,103,104,105,106,107,108,109,110,111,112,113,114,115,116,117,118].

# 2. Applications in Economics and Finance

## 2.1. Prices and Volumes

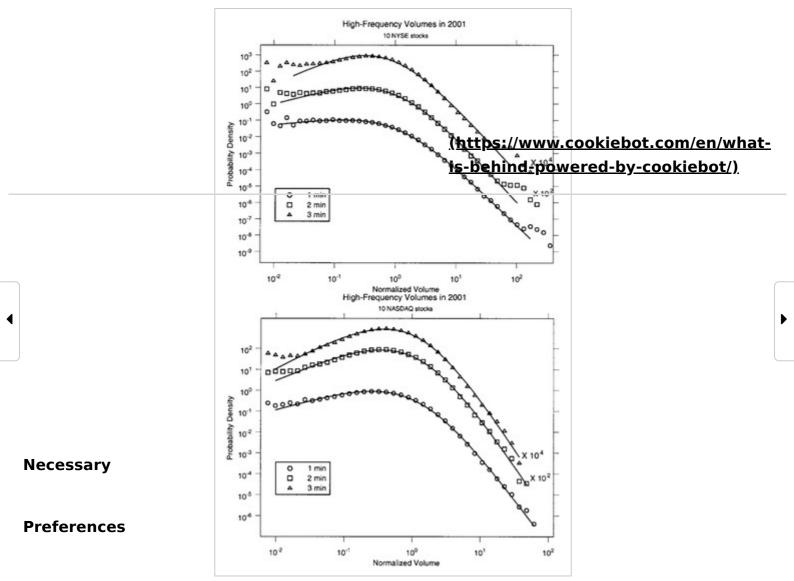
Time series of prices (say of stocks, commodities, etc.), where t runs along chosen units (say seconds or minutes, or days, or years) are conveniently replaced by their corresponding returns (or logarithmic returns), defined as follows:

Returns do not depend on the specific currency of the prices and fluctuate around zero; in addition to that, their definition cancels systematic inflation. The distribution of returns usefully characterizes the price fluctuations. See an illustration in **Figure 1**, from [103] (see also [104,118]). The amounts of the corresponding transactions are currently referred to as volumes: see, for example, **Figure 2**.



**Figure 1.** Empirical return densities (points) and q-Gaussians (solid lines) for normalized returns of the 10 top-volume stocks in the NYSE and in NASDAQ in 2001. The dotted line is a (visibly inadequate) Gaussian distribution. The 2- and 3-min curves are moved vertically for display purposes. From [103]. There exist in the literature quite a few other such examples, for other stocks and other years, with similar values of q.

Show details



#### **Statistics**

**Figure 2.** Empirical distributions (points) and *q*-exponential-like curves (solid lines) for normalized volumes of the 10 top-volume stocks in the NYSE and in the NASDAQ in 2001. The **Marketing** solid lines are fittings with a *q*-exponential multiplied by a power-law (analogous to the *u*-ensity of states prefactor that typically emerges for the distributions of quasi-particles in, say, condensed matter physics); from [103]. There exist in the literature quite a few other such examples, for other stocks and other years, with similar values of *q* and of the rest of the fitting indices.

#### 2.2. Volatilities

The volatility characterizes the size (standard deviation) of the fluctuations of returns. The volatility smile characterizes the correction of empirical volatilities with regard to a Gaussian-based expectation: see an illustration in **Figure 3** (from [103]). To be more explicit, let us assume that we are handling the following Gaussian distribution, where characterizes univocally the volatility. To discuss the probability distribution of quantities such as, Queiros introduced [108] the *q*-log normal probability function:

(21)

where is a normalizing factor and are parameters. The particular case corresponds to the standard log-normal function. See **Figure 4** for illustrative examples of this function. See also **Figure 5** for a real financial application.

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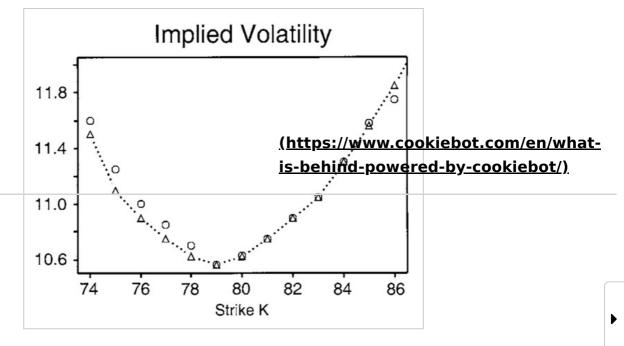
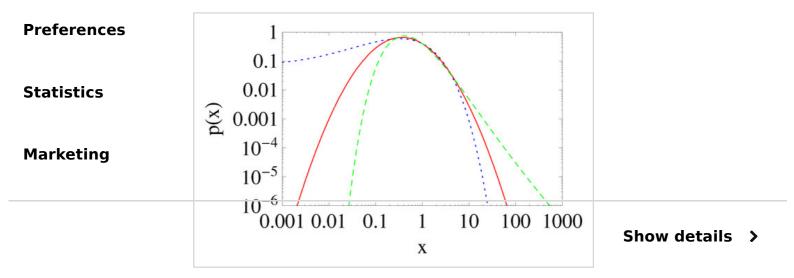
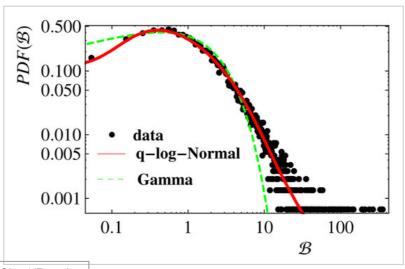


Figure 3. Implied volatilities as a function of the strike price for call options on JY currency futures, traded on 16 May 2002, with 147 days left to expiration. In this typical example, the current price of a contract on Japanese futures is \$79, and the risk-free rate of return is . Circles correspond to volatilities implied by the market, whereas triangles correspond to volatilities NARGARD our model with and . The dotted line is a guide to the eye. From [103].



**Figure 4.** Illustrations of the q-log-normal density for and : blue , red and green . From [108].

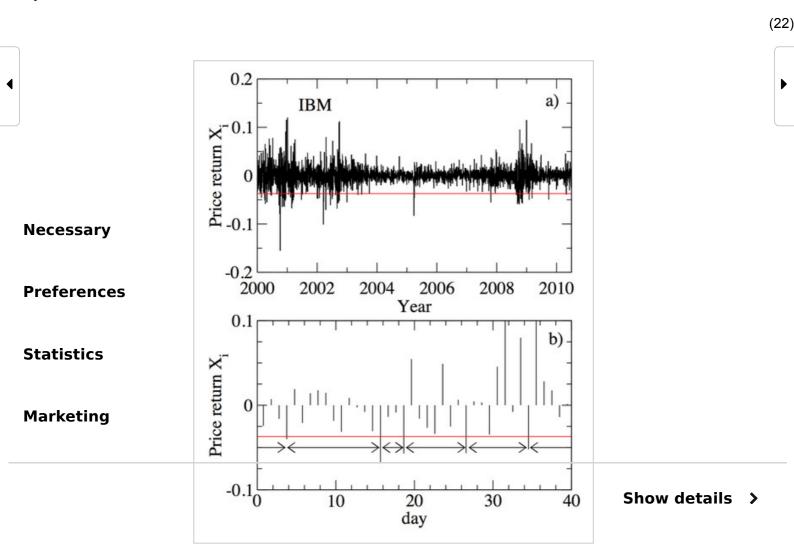


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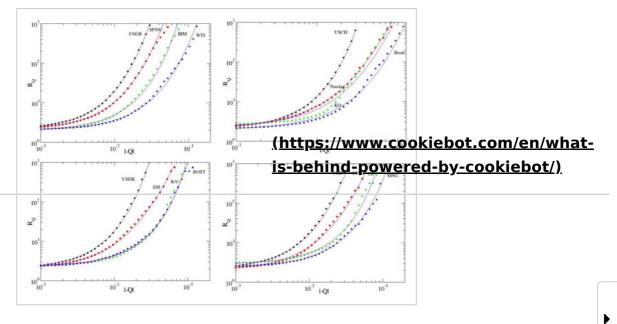
**Figure 5.** Probability density function of a five-day volatility vs. . The symbols are obtained from the data, and the lines are the best fits with the Gamma distribution (dashed green) and the double-sided q-log-normal (red) with , and . For further details, see [108].

### 2.3. Inter-Occurrence Times

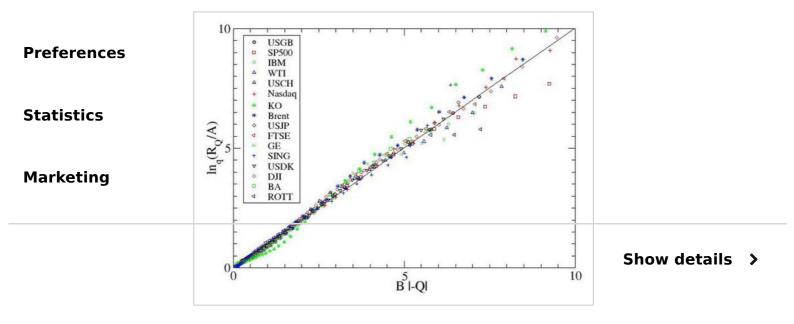
We can see in **Figure 6** two typical time series of price the price the price the price of the corresponds to an average inter-occurrence time to the examples shown in **Figure 7** and **Figure 8**. For a fixed value of , we verify that , with:



**Figure 6.** Illustration of the relative daily price returns of the IBM stock between (**a**) January 2000 and June 2010 and (**b**) 27 August and 23 October 2002. The red line shows the threshold , which corresponds to an average inter-occurrence time of = 70. In (b), the inter-occurrence times are indicated by arrows. From [106].



**Figure 7.** The mean inter-occurrence time vs. the absolute value of the loss threshold –Q. The continuous curves are fittings with . Top left: For the exchange rate of the U.S. Dollar against the British Pound, the index S&P500, the IBM stock and crude oil (West Texas Intermediate (WTI)), from left to right in the plot; the corresponding values for are 0.95, 0.92, 0.97, 0.927 (with and ). Similarly for the top right, bottom left and bottom right plots. From [107]. **Necessary** 

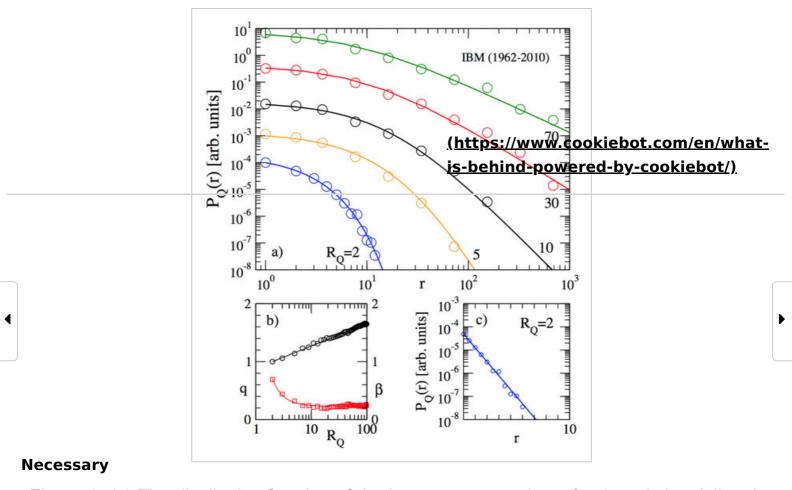


**Figure 8.** The mean inter-occurrence time versus the absolute value of the loss threshold  $\neg Q$ : versus the representation of the same data of **Figure 7**. The continuous curve is a fitting with . From [107].

See the illustrations in **Figure 9** and **Figure 10**. The fact that we have analytically enables us to straightforwardly obtain an explicit expression for the risk function, which is defined as the probability of having once again a fluctuation larger than within an interval at time *t* after the last large fluctuation. It can be shown [106,109] that:

(23)

with . See **Figure 11**.



**Figure 9.** (a) The distribution function of the inter-occurrence times for the relative daily price **Pretrum note** M in the period 1962–2010. The data points belong to = 2, 5, 10, 30 anc 0 (in units of days), from bottom to top. The full lines show the fitted *q*-exponentials for typical values of . (b) The dependence of the parameters (squares, lower curve) and (circles, upper curve) on **Statistics** in the -exponential. (c) Confirmation that, for , the distribution function is a simple exponential

(i.e., ). The straight line is proportional to . From [106].

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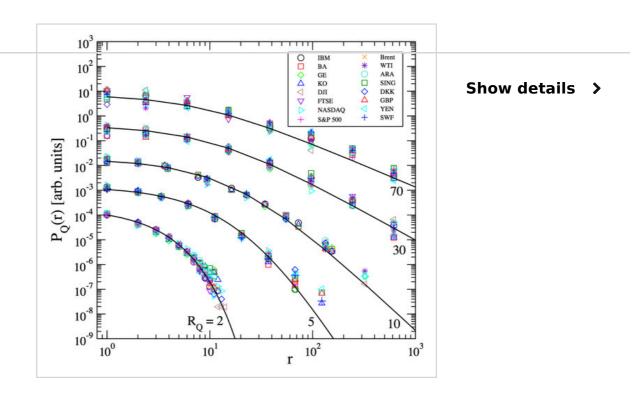
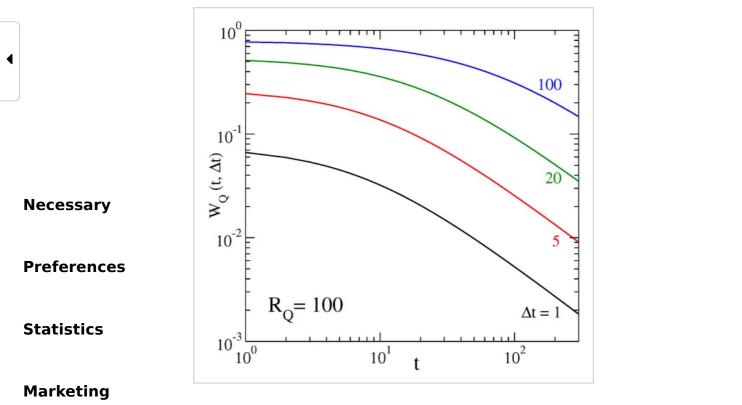


Figure 10. The distribution function of the inter-occurrence times (as in Figure 9a) for the relative daily price returns of 16 examples of financial data, taken from different asset classes (stocks, indices, currencies, commodities). The assets are: (i) the stocks of IBM, Boeing (BA), General Electric (GE), Coca-Cola (KO); (ii) the indices Dow Jones (DJI), Financial Times Stock Exchange 100 (FTSE), NASDAQ, S&P 500; (iii) the commodities Brent Crude Oil, West Texas Intermediate (WTI), Amsterdam-Rotterdam-Antwerphysical (Indiana Personal Personal



**Figure 11.** Universal risk function from Equation (**23**) for the inter-occurrence time = 100 and for the intervals = 1, 5, 20, 100 days (from bottom to top). From [**106**].

### 2.4. Wealth

Show details >

Wealth inequality within a given country is a classical and most important matter, which can be characterized within *q*-statistics as shown in [105]: see Figure 12 and Figure 13. The larger the index is, the larger the inequality. As we verify, the U.K. and Germany are more egalitarian countries than the U.S. and Brazil. In addition to that, inequality appears to increase in the U.S. and Brazil, at least during the years indicated in the plots.

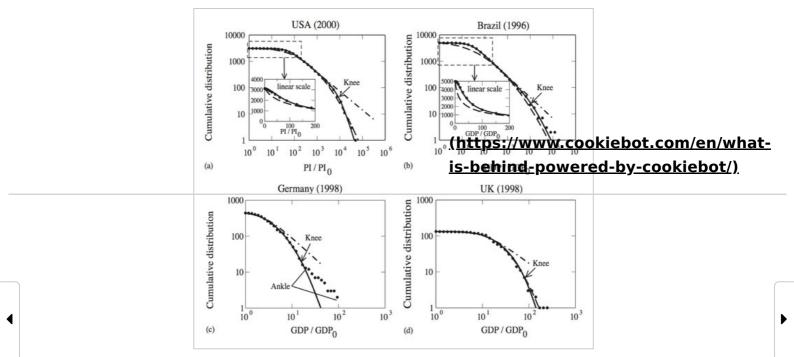
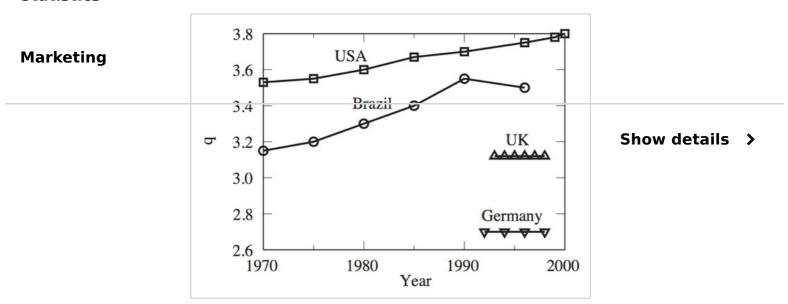


Figure 12. Binned inverse cumulative distribution of the county, (U.S.) and (Brazil, Germany and U.K.), where denotes the Personal Income and denotes the Gross Domestic Product of countries. Three distributions are displayed for comparison: (i) *q*-Gaussian (with) (dot-dashed); Nieces and use (solid) and (iii) log-normal (dashed lines). (a,b) present insets with a lin relinear scale, to make more evident the quality of the fitting at the low region (in (c,d), the ()-Gaussian and the log-normal curves are superposed and, so, are visually indistinguishable). Preferences

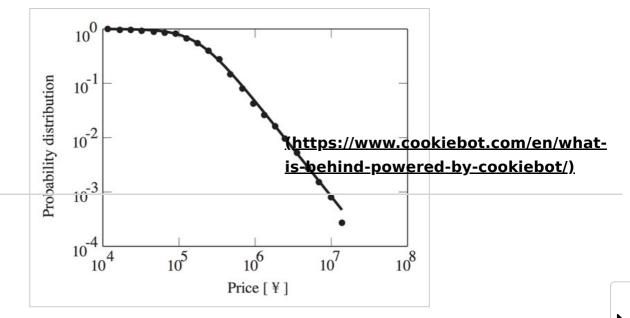
The positions of the knees are indicated. The ankle is particularly pronounced in (c), though it is also present in the other cases. From [105], where further details are available.

#### **Statistics**



**Figure 13.** Evolution of parameter q for the U.S. (squares), Brazil (circles), the U.K. (up triangles) and Germany (down triangles). The parameters (for each country) are constant for all years: = 2.1, = 1.7, = 1.5, = 1.4. Lines are only guides to the eyes. As we verify, in some cases, the index q remains invariant along time, whereas in others, it evolves; the functional forms remain however the same as indicated in **Figure 12**. From [105].

Another index that characterizes the wealth of a country and its inequalities is associated with the prices of the land. See in Figure 14 (from [105]) an illustration for Japan, where .



**Figure 14.** Inverse cumulative probability distribution of Japanese land prices for the year 1998. The solid curve is a q-Gaussian with , which corresponds to the slope , and 188,982 Yen. From [105], where further details are available.

We have described a variety of financial and economic properties with a plethora of q-indices,

#### 3. Conclusions and Perspectives Necessary

such as. For a given system, how many independent indices should we expect? The full answer to this question remains up to now elusive. It seems however that only a few of them are essentially independent, all of the others being (possibly simple) functions of those few. Such an algebraic structures first advanced and described in [119] and has been successfully verified in a solar wind [53] (see also [6] and the references therein) and elsewhere; it has recently been generalized [120,121] and related to the Moebius group. The central elements of these algebraic and possible verification of such structures constitutes nowadays an important open question, whose further study would surely be most useful.

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Another crucial question concerns the analytic calculation from first principles of some or all of the above *q*-indices. This is in principle possible (as illustrated in [63,64,65,66]), but it demands the complete knowledge of the microscopic model of the specific class of the complex system. For the full set of the *q*-indices shown in the present overview, such models are not available, even if they would be very welcome.

Let us finally emphasize that many other statistical approaches exist for the quantities focused on in the present overview. However, as announced in the title of this paper, this is out of the present scope. The present paper is one among various others belonging to the same Special Issue of the journal Entropy. The entire set of articles is expected to enable comparisons between these many approaches.

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especially J. Ludescher for authorizing me to use, in the present review, our unpublished figures of [1]

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Conflicts of Interest

thermodynamics.

The author declares no conflict of interest.

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